

QCD24

27th HIGH-ENERGY PHYSICS
INTERNATIONAL CONFERENCE
IN QUANTUM CHROMODYNAMICS



$a_\mu|_{l.o}^{hvp}$, **power corrections and α_s from e^+e^-**

Stephan Narison



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 - QCD condensates are used as inputs in some other approaches.

Based on the recent papers

- ♣ QCD parameters and SM-high precision from $e^+e^- \rightarrow$ Hadrons:
Updated: [SN](#), [Nucl.Phys.A 1046 \(2024\) 122873](#) : [ArXiv 2402.13983 \[hep-ph\]](#)

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[Nucl.Phys.A 1039 \(2024\) 122744](#) : [ArXiv 2306.14639 \[hep-ph\]](#)

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- ♥ **Summary:** α_s from e^+e^- and τ -like decay BNP moment



$e^+e^- \rightarrow \text{Hadrons}$
Total cross-sections
 $R = \text{Hadrons} / \mu^+\mu^-$

History : Dispersion Relation

- Bridge between High AND Low energy QCD regions

$$\begin{aligned}\Pi_H(Q^2 \equiv -q^2) &\equiv i \int d^4x \langle 0 | \mathcal{T} J_H(x) J_H^\dagger(0) | 0 \rangle \\ &= \int_{t < \infty} \frac{dt}{t + Q^2 + i\epsilon} \text{Im}\Pi(t) + \text{subtraction terms}\end{aligned}$$

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- $J_H^\mu(x) = \sum_q e_q^2 \bar{\Psi}_q \gamma^\mu \Psi_q$: electromagnetic current

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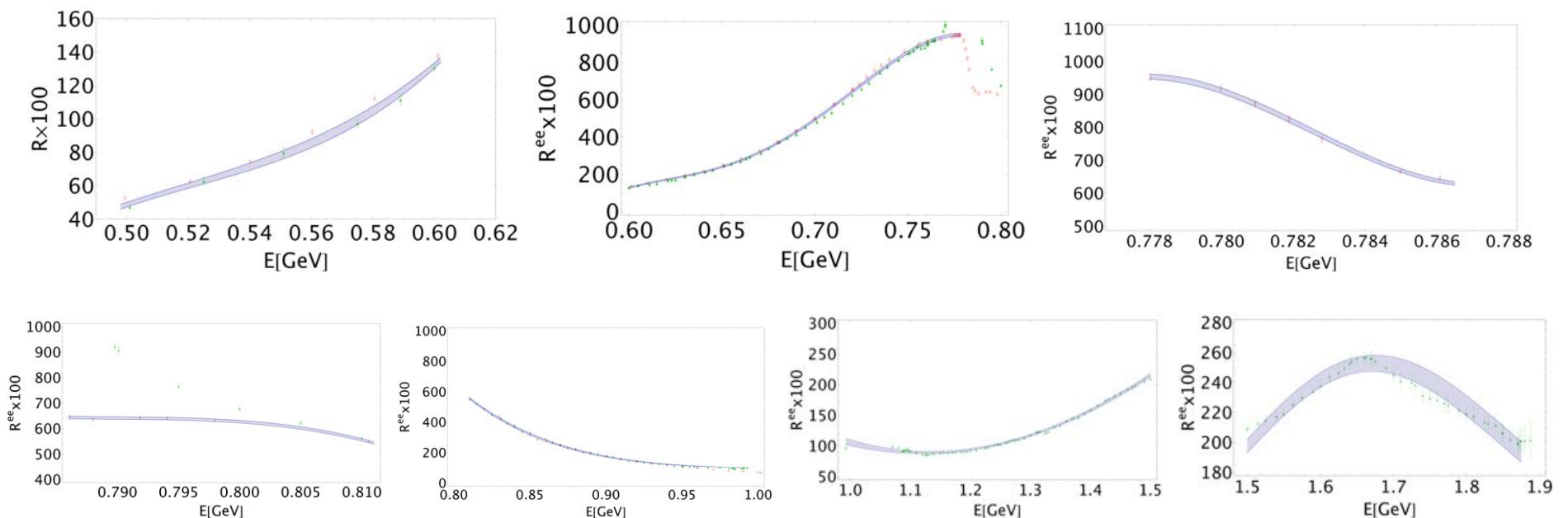
- $J_H^\mu(x) = \sum_q e_q^2 \bar{\Psi}_q \gamma^\mu \Psi_q$: electromagnetic current

- $\frac{1}{\pi} \text{Im}\Pi(t) \sim \sigma_{tot}(e^+e^- \rightarrow \text{Hadrons})$: Optical theorem.

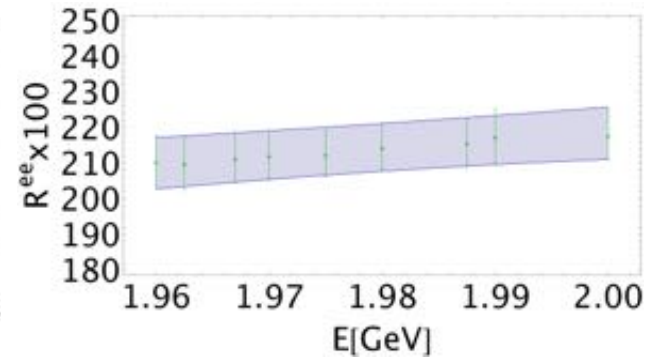
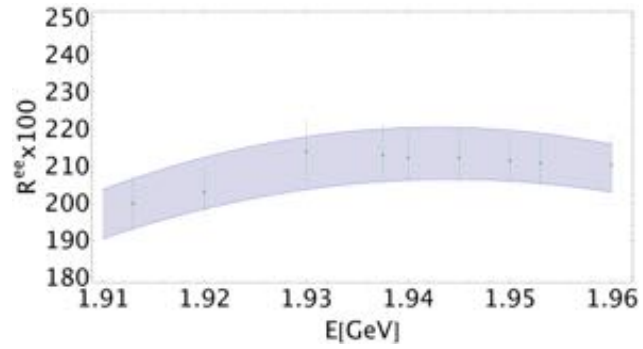
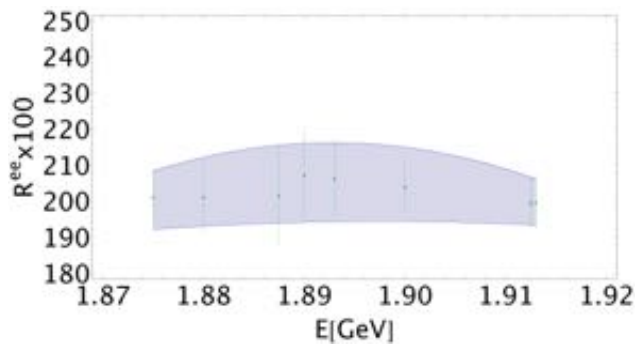
• ♣ Data compiled by PDG22 ⊕ New CMD3 ($E \leq 1.875$ GeV)

$$R^{ee} \equiv \frac{\sigma(e^+e^- \rightarrow \text{Hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} = 12\pi \text{Im}\Pi(t) : \text{ Optical theorem}$$

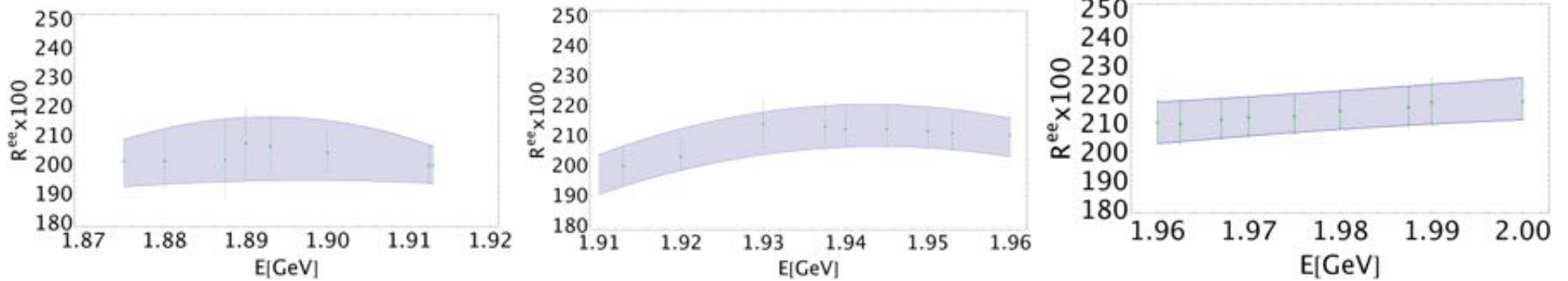
We divide the region $0.5 \leq \sqrt{t} \leq 1.9$ GeV into 7 subregions and use a Mathematica Interpolating Fitting program.



● Data fit from 1.875 to 2 GeV



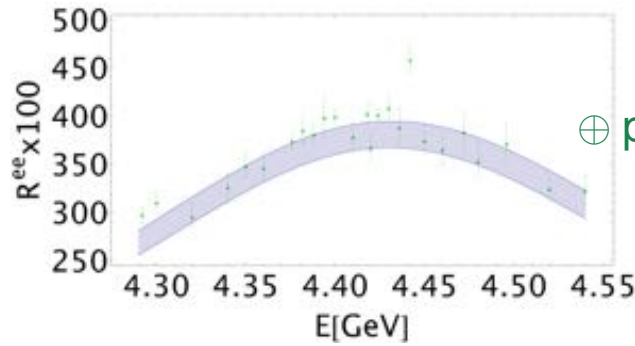
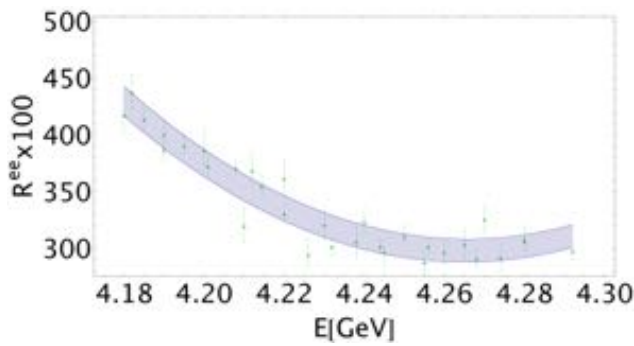
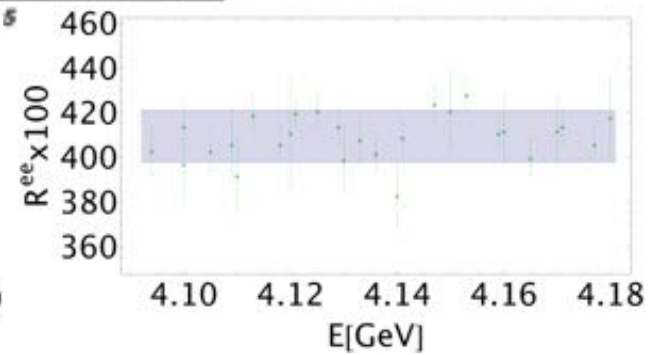
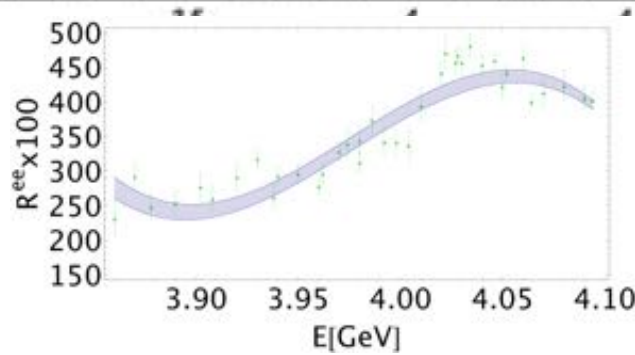
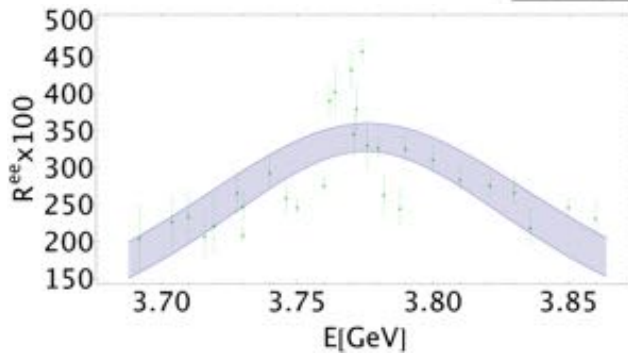
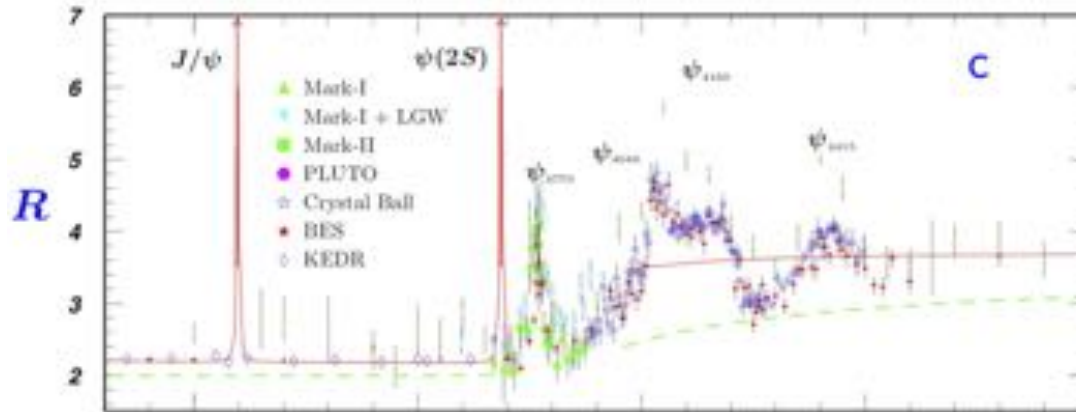
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- pQCD - 3 flavours from 2 to 3.68 GeV

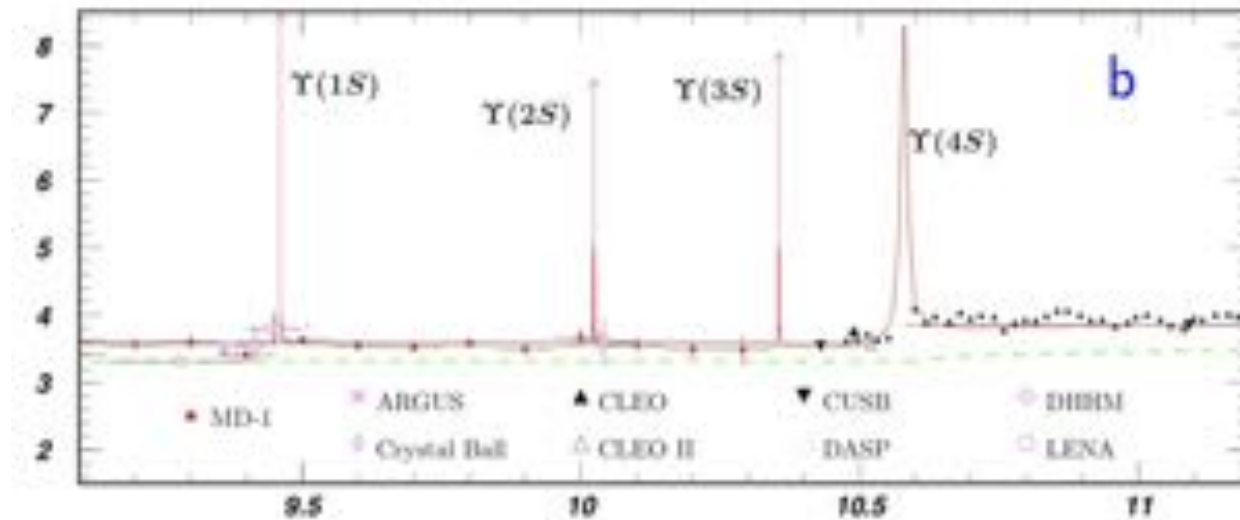
$$\begin{aligned}
 R &= 3 \sum Q_i^2 \left[1 + a_s + 1.6398 a_s^2 - 10.2839 a_s^3 - 106.8798 a_s^4 + O(a_s^5) \right. \\
 &+ d_{2m} = -\frac{\bar{m}^2}{t} (12 a_s + 104.833 a_s^2 + 541.735 a_s^3), \\
 &+ d_{4m} = \frac{\bar{m}^4}{t^2} \left[-6 \left(1 + \frac{11}{3} a_s \right) + a_s^2 \left[n_f \left(\frac{L_m}{3} - 1.841 \right) - \frac{11}{2} L_m \right. \right. \\
 &+ \left. \left. 136.693 + 12 - 0.475 - L_m \right] \right] : L_m = \text{Log}(\bar{m}^2 / E^2)
 \end{aligned}$$

● Detailed fit of charmonium data from 3.68 to 4.55 GeV



⊕ pQCD - 4 flavours ($E \geq 4.58$ GeV)

● Bottomium data from PDG 22



$\Upsilon(1S), \dots, \Upsilon(4S)$: Narrow Width Approximation

⊕ pQCD - 5 flavours ($E \geq 10.6$ GeV)



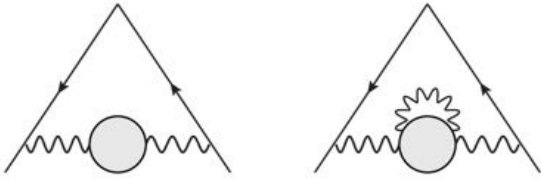
TEST OF THE SM

from $(g-2)_\mu$

using $e^+e^- \rightarrow \text{Hadrons}$

Test / Calibration of the Fit and Estimate of $a_\mu|_{l.o}^{hvp}$

- Hadronic Vacuum Polarisation of $a_\mu|_{l.o}^{hvp} \equiv \frac{1}{2}(g-2)_\mu|_{l.o}^{hvp}$

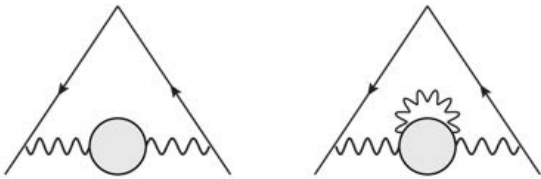


$\frac{1}{\pi} \text{Im}\Pi(t)|_{l.o}^{hvp} \sim \sigma_{tot}(e^+e^- \rightarrow \text{Hadrons})$: Optical theorem.

$$a_\mu|_{l.o}^{hvp} = \frac{1}{4\pi^3} \int_{4m_\pi^2}^{\infty} dt K_\mu(t/m_\mu^2) \sigma(e^+e^- \rightarrow \text{hadrons})$$

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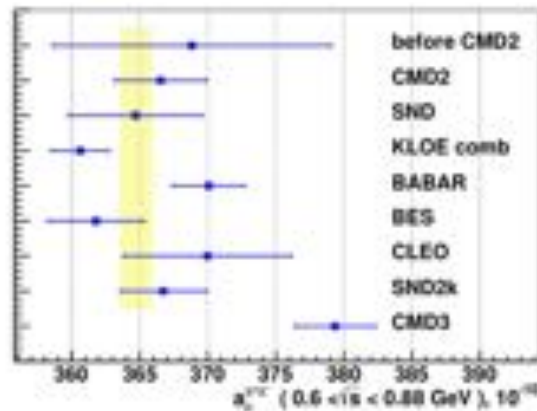
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- Comparison of \neq Results in the Low-Energy regions



Experiment	$a_\mu^{\rho^+ \pi^-, LO}, 10^{-10}$
before CMD2	368.8 ± 10.3
CMD2	366.5 ± 3.4
SND	364.7 ± 4.9
KLOE	360.6 ± 2.1
BABAR	370.1 ± 2.7
BES	361.8 ± 3.6
CLEO	370.0 ± 6.2
SND2k	366.7 ± 3.2
CMD3	379.3 ± 3.0

$$a_\mu^{\rho}|_{l.o.}^{hvp} [0.6 \rightarrow 0.88] = (377.4 \pm 3.1) \times 10^{-10}, \quad a_\mu^{\rho}|_{l.o.}^{hvp} [0.6 \rightarrow 0.993] = (400.9 \pm 3.2) \times 10^{-10}.$$

About the same as CMD3.

$$\implies a_\mu^{I=1+0}|_{l.o.}^{hvp} [2m_\pi \rightarrow 1.875] = (6492.3 \pm 38.8) \times 10^{-11},$$

(100 ~ 136) $\times 10^{-11}$ larger than Davier et al., Nomura et al. using KLOE data.

Final results

$$a_{\mu|lo}^{hvp} = (7036.5 \pm 36.9) \times 10^{-11}$$

SN1978 (3rd cycle thesis) :

$$(7020 \pm 800) \times 10^{-11}$$

Same central value but

22 × more precise ! ⇒

$$a_{\mu}^{exp} - a_{\mu}^{th} = (142 \pm 42_{th} \pm 22_{exp}) \times 10^{-11}$$

3σ deviation from SM : New BNL data 2023

$$a_{\tau}^{hvp} = (3494.8 \pm 24.7) \times 10^{-9}$$

SN 1978 : $(3700 \pm 400) \times 10^{-9}$: 1st estimate !

$$\Delta\alpha_{had}^{(5)}(M_Z) = (2766.3 \pm 4.5) \times 10^{-5}$$

\sqrt{s} [GeV]	$a_{\mu lo}^{hvp} \times 10^{11}$	$a_{\tau}^{hvp} \times 10^9$	$\Delta\alpha_{had}^{(5)}(M_Z^2) \times 10^5$
Light I=1			
$\rho(2m_{\pi} \rightarrow 0.50)$	489.4 ± 3.1	85.1 ± 0.6	12.38 ± 0.08
$\rho(0.50 \rightarrow 0.60)$	524.5 ± 13.7	135.9 ± 3.5	22.84 ± 0.59
$\rho(0.60 \rightarrow 0.776)$	2712.0 ± 30.2	943.7 ± 10.3	182.87 ± 1.97
$\rho(0.776 \rightarrow 0.993)$	1297.3 ± 9.8	1797.3 ± 20.3	117.82 ± 0.98
0.993 → 1.5	354.4 ± 6.7	228.2 ± 4.1	67.3 ± 1.14
1.5 → 1.875	237.6 ± 5.7	206.8 ± 4.9	80.11 ± 1.9
Total Light I=1 (≤ 1.875)	5615.2 ± 36.0	2452.9 ± 11.9	489.9 ± 4.4
Light I=0			
ω (NWA)	417.1 ± 13.7	163.3 ± 5.3	31.6 ± 1.1
ϕ (NWA)	380.6 ± 4.6	20.6 ± 0.2	51.2 ± 0.6
0.993 → 1.5	44.3 ± 0.8	28.5 ± 0.5	8.4 ± 0.1
$\omega(1650)$ (BW)	24.3 ± 0.1	16.7 ± 0.1	5.2 ± 0.1
$\phi(1680)$ (BW)	1.8 ± 0.9	1.3 ± 0.6	0.4 ± 0.2
Total Light I=0 (≤ 1.875)	877.7 ± 14.5	230.4 ± 5.7	99.4 ± 1.7
Light I=0 ⊕ 1			
1.875 → 1.913	14.7 ± 0.7	14.4 ± 0.7	6.3 ± 0.3
1.913 → 1.96	17.6 ± 0.6	17.5 ± 0.6	7.9 ± 0.3
1.96 → 2	14.3 ± 0.5	14.5 ± 0.5	6.7 ± 0.2
2 → 3.68: QCD (u, d, s)	247.2 ± 0.3	308.3 ± 0.5	202.8 ± 0.5
Total Light I=0 ⊕ 1 (1.875 → 3.68)	297.8 ± 1.1	354.7 ± 1.2	223.7 ± 0.7
Total Light I=0 ⊕ 1 (2m_π → 3.68)	6796.1 ± 38.8	3036.4 ± 24.5	806.4 ± 4.6
Charmonium			
$J/\psi(1S)$ (NWA)	65.1 ± 1.2	92.7 ± 1.8	73.5 ± 1.4
$\psi(2S)$ (NWA)	16.4 ± 0.6	26.0 ± 0.9	26.1 ± 0.8
$\psi(3773)$ (NWA)	1.7 ± 0.1	2.7 ± 0.2	2.9 ± 0.2
Total J/ψ (NWA)	83.2 ± 1.4	121.4 ± 2.9	102.5 ± 1.6
3.69 → 3.86	11.4 ± 1.0	18.3 ± 1.6	19.0 ± 1.6
3.86 → 4.094	16.6 ± 0.5	27.5 ± 0.8	30.9 ± 0.9
4.094 → 4.18	6.6 ± 0.2	11.2 ± 0.3	13.2 ± 0.4
4.18 → 4.292	6.5 ± 0.2	11.2 ± 0.4	13.7 ± 0.5
4.292 → 4.54	11.8 ± 0.6	20.7 ± 0.8	26.8 ± 1.1
4.54 → 10.50: QCD (u, d, s, c)	92.0 ± 0.0	186.2 ± 0.0	458.73 ± 0.1
Total Charmonium (3.69 → 10.50)	145.7 ± 1.7	275.6 ± 2.9	364.9 ± 2.2
Total Charmonium	228.1 ± 1.9	396.5 ± 2.8	604.8 ± 2.7
Bottomium			
$\Upsilon(1S)$ (NWA)	0.54 ± 0.02	1.25 ± 0.07	5.65 ± 0.29
$\Upsilon(2S)$ (NWA)	0.22 ± 0.02	0.51 ± 0.06	2.54 ± 0.29
$\Upsilon(3S)$ (NWA)	0.14 ± 0.02	0.33 ± 0.04	1.77 ± 0.21
$\Upsilon(4S)$ (NWA)	0.10 ± 0.01	0.23 ± 0.03	1.26 ± 0.16
$\Upsilon(10.86 \oplus 11)$ (NWA)	0.1 ± 0.06	0.20 ± 0.06	1.67 ± 0.39
Total Bottomium (NWA)	1.0 ± 0.1	2.3 ± 0.1	11.3 ± 0.5
Z - pole	-	-	-29.2[8]
10.59 → 2m _τ : QCD (u, d, s, c, b)	22.4 ± 0.3	57.5 ± 0.1	1282.9 ± 1.2
Total Bottomium	23.4 ± 0.3	59.8 ± 0.1	1323.3 ± 1.3
2m _τ → ∞: QCD (u, d, s, c, b, t)	0.03	0.08	-28.2
Total sum	7036.5 ± 38.9	3494.8 ± 24.7	2766.3 ± 4.5

Comparisons with Results reported @ QCD 24

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- BABAR (G. Vasseur)

$$\begin{aligned} a_\mu^{exp} - a_\mu^{th} &= (168 \pm 38 \pm 29) : \text{BABAR} \\ &= (142 \pm 42_{th} \pm 22_{exp}) : \text{This Work } (3\sigma) \end{aligned}$$

◇ *SVZ sum rule* 45 years ago !



Zakharov-fest Ringberg Castel, Tegernsee, Munich (19-21 May 2005).
From left to right: A.I. Vainshtein, Your Servitor, V.I. Zakharov, M.A. Shifman.

History : Dispersion Relation

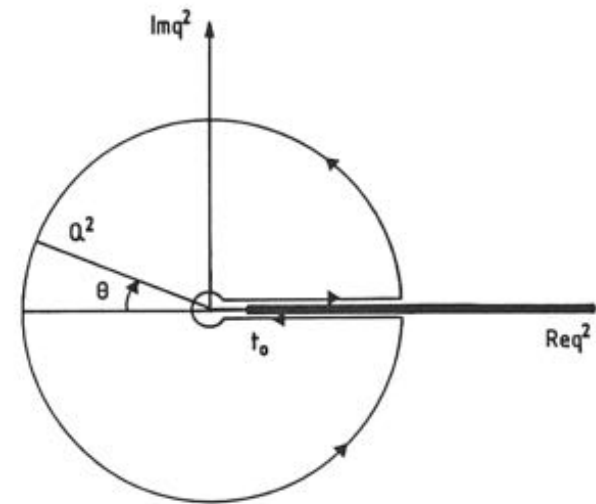
- Bridge between High AND Low energy QCD regions

$$\begin{aligned}\Pi_H(Q^2 \equiv -q^2) &\equiv i \int d^4x \langle 0 | \mathcal{T} J_H(x) J_H^\dagger(0) | 0 \rangle \\ &= \int_{t < \infty} \frac{dt}{t + Q^2 + i\epsilon} \text{Im}\Pi(t) + \text{subtraction terms}\end{aligned}$$

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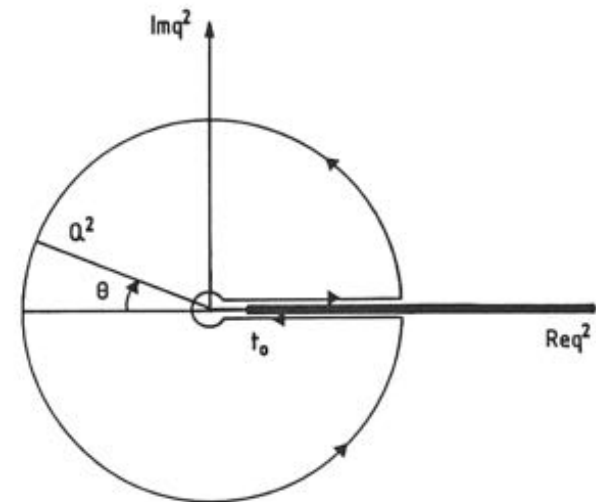


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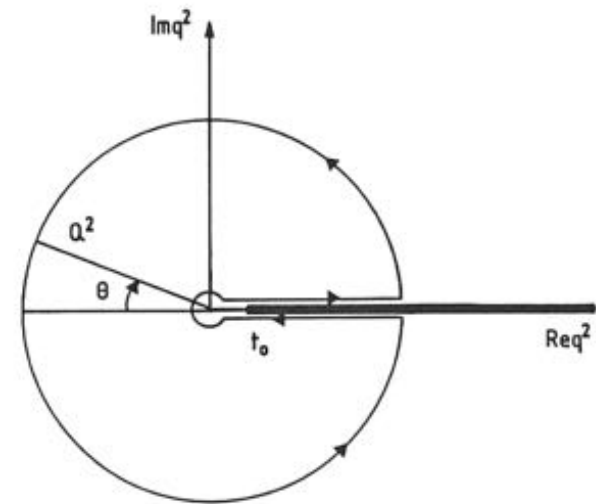
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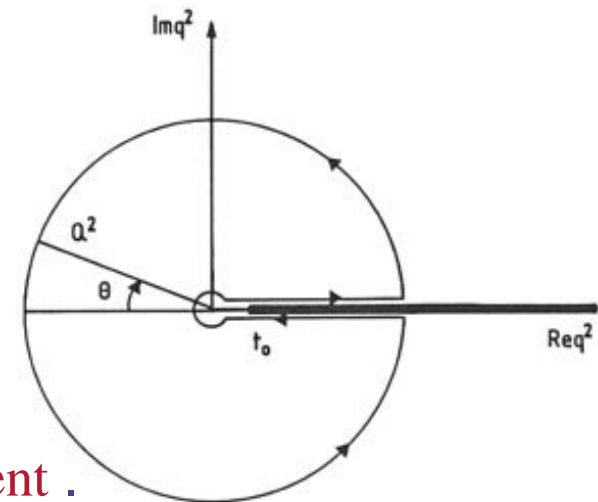
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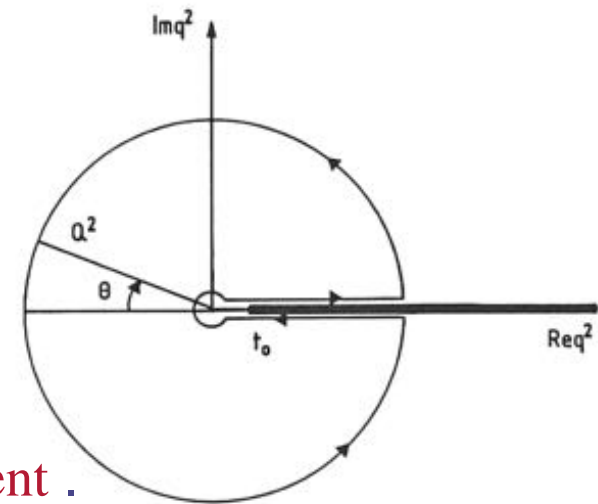
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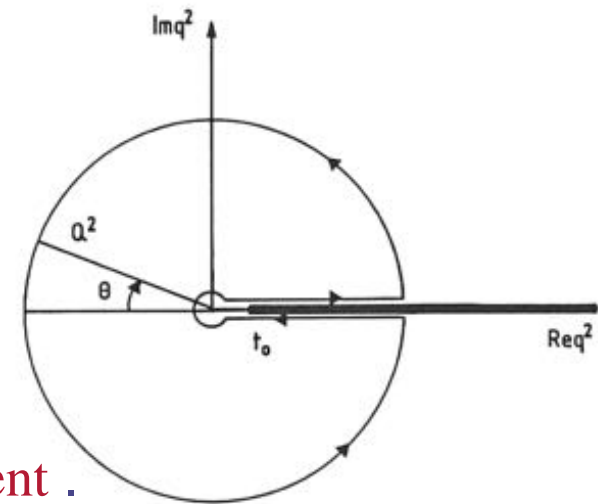
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- ♣ Example of Vector two-point function

$$8\pi^2\Pi(Q^2) = \sum_{D=0,2,4,\dots} \frac{C_D \langle 0|O_D|0\rangle}{Q^D} : \quad d_D \equiv C_D \langle 0|O_D|0\rangle$$

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♡ Truncation of the OPE

- A truncation of the OPE up to $D = 6$ is enough for Phenomenology !
- No good control of condensates beyond $D = 6$: violation of factorization, mixing under renormalization NT83, often some classes of diagrams are only computed,...

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- ♣ Improvement of the dispersion relation

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$$\mathcal{R}(\tau) \equiv -\frac{d}{d\tau} \log \mathcal{L}(\tau) \simeq M_R^2, \quad r_{12}(\tau_H) \equiv \frac{\mathcal{R}_1}{\mathcal{R}_2} \simeq \frac{M_{R1}^2}{M_{R2}^2}$$

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- $\mathcal{R}(\tau)$: less sensitive to α_s -corrections \implies
Good tools for extracting the QCD condensates

Launer-SN-Tarrach 84



QCD CONDENSATES
from
 $e^+e^- \rightarrow$ Hadrons

QCD condensates from Ratio of LSR

- Choice of the ratio of LSR **Launer-SN-Tarrach 84**

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- $\mathcal{R}_{10}(\tau)$ less sensitive to α_s than the moment $\mathcal{L}_0(\tau)$: starts to $O(\alpha_s^2)$.

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$d_6 = -\frac{896}{81} \pi^3 \rho \alpha_s \langle \bar{\Psi}_q \Psi_q \rangle^2$: four-quark condensates : $\rho = 1$: factorization.

$d_8 = \langle GGGG \rangle$: 4-gluon condensate $\oplus \dots$

- $\mathcal{R}_{10}(\tau)$ less sensitive to α_s than the moment $\mathcal{L}_0(\tau)$: starts to $O(\alpha_s^2)$.
- Relative contributions of the condensates increased in the OPE !

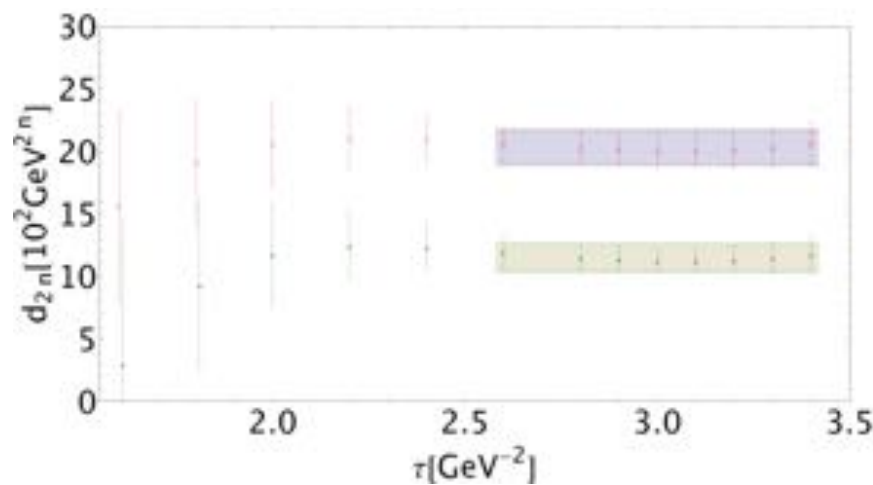
Fitting procedure

- 3-parameter fit (d_4, d_6, d_8) : not conclusive ! Like τ -decay Moments !

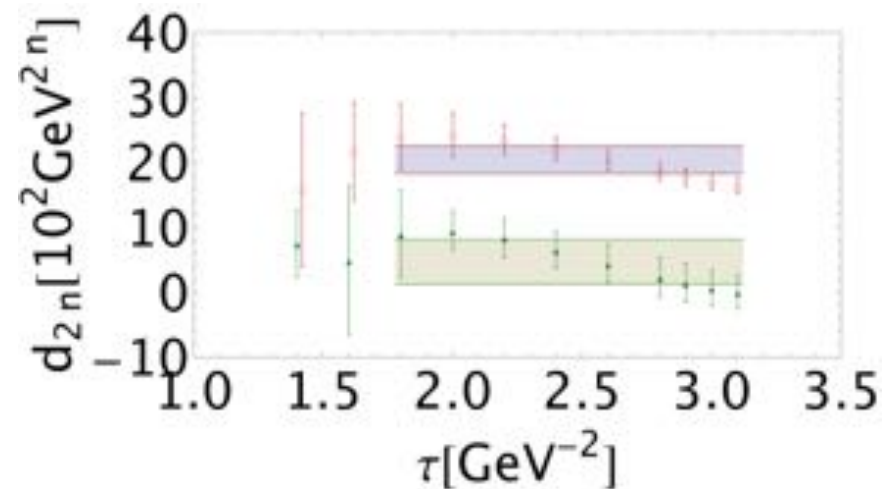
Fitting procedure

- 3-parameter fit (d_4, d_6, d_8) : **not conclusive !** Like τ -decay Moments !
- $\langle \alpha_s G^2 \rangle$ from heavy quarkonia \oplus 2-parameter fit of (d_6, d_8) :

$O(\alpha_s^2)$



$O(\alpha_s^4)$



d_6 stable for $\alpha_s^2 \rightarrow \alpha_s^4$ but not d_8 . To order α_s^4 :

$$d_6 = -(20.5 \pm 2.0) \times 10^{-2} \text{ GeV}^6, \quad d_8 = (4.7 \pm 3.5) \times 10^{-2} \text{ GeV}^8.$$

NB : $\tau \simeq (2 \sim 3) \text{ GeV}^{-2}$ relatively big ! Results to be improved later on !

τ -like decay moments BNP 92

$$R_n^{ee} = \int_0^1 dx_0 (1 - 3x_0^2 + 2x_0^3) x_0^n 2R_{ee}^{I=1}(x_0) : x_0 \equiv (t/M_0^2)$$

- Perturbative corrections $\delta_n^{(0)}$ to order α_s^4 (FO) for $n \geq 1$, :

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- Perturbative corrections $\delta_n^{(0)}$ to order α_s^4 (FO) for $n \geq 1$, :
- Condensates to lowest order of α_s **Cauchy Property** :

$$R_0^{ee} = f(\delta_0^{(0)}, d_6, d_8), \quad d_{n \geq 10} = 0$$

$$R_1^{ee} = f(\delta_1^{(0)}, d_4, d_8, d_{10}), \quad d_{n \geq 12} = 0$$

$$R_2^{ee} = f(\delta_2^{(0)}, d_6, d_{10}, d_{12}), \quad d_{n \geq 14} = 0$$

$$R_3^{ee} = f(\delta_3^{(0)}, d_8, d_{12}, d_{14}), \quad d_{n \geq 16} = 0$$

$$R_4^{ee} = f(\delta_4^{(0)}, d_{10}, d_{14}, d_{16}), \quad d_{n \geq 18} = 0$$

$$R_5^{ee} = f(\delta_5^{(0)}, d_{12}, d_{16}, d_{18}), \quad d_{n \geq 20} = 0$$

$$R_6^{ee} = f(\delta_6^{(0)}, d_{14}, d_{18}, d_{20}), \quad d_{n \geq 22} = 0$$

QCD condensates from τ -like decay moments

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- Initial Inputs : α_s , and $\langle \alpha_s G^2 \rangle$ ($J/\psi, \Upsilon$) , d_6 (LSR) condensates \implies
 R_0 gives $d_8 \implies R_1$ gives $d_{10} \implies R_2$ gives $d_{12} \implies R_3$ gives
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QCD condensates from τ -like decay moments

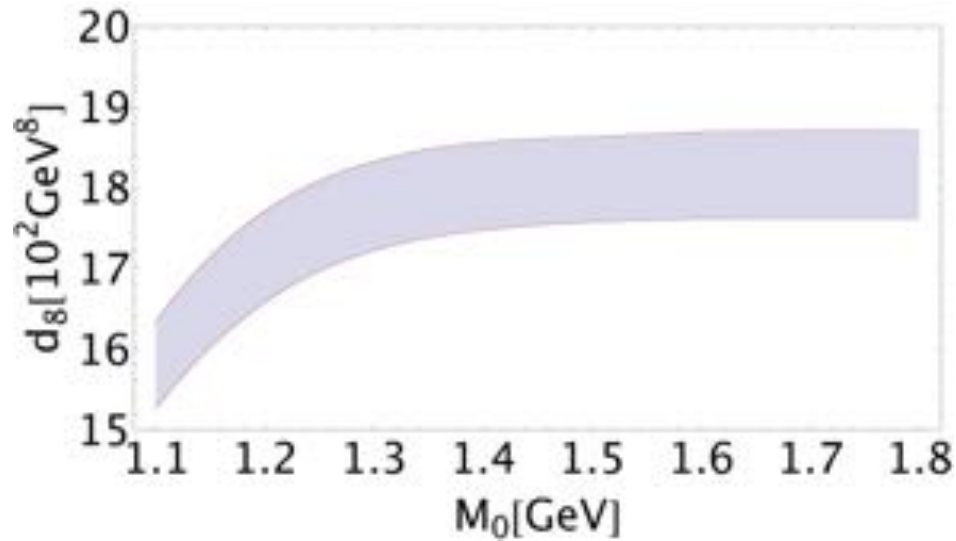
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- We proceed iteratively to improve the results

$$\begin{array}{cccccccccccc}
 d_4 & \longleftrightarrow & d_6 & \longleftrightarrow & d_8 & \overset{(2)}{\longleftrightarrow} & d_{10} & \overset{(6)}{\longleftrightarrow} & d_{12} & \longrightarrow & d_{14} & \longrightarrow & d_{16} \\
 J/\psi, \Upsilon & LSR & & R_0 & & R_1 & & R_2 & & R_3 & & R_4 &
 \end{array}$$

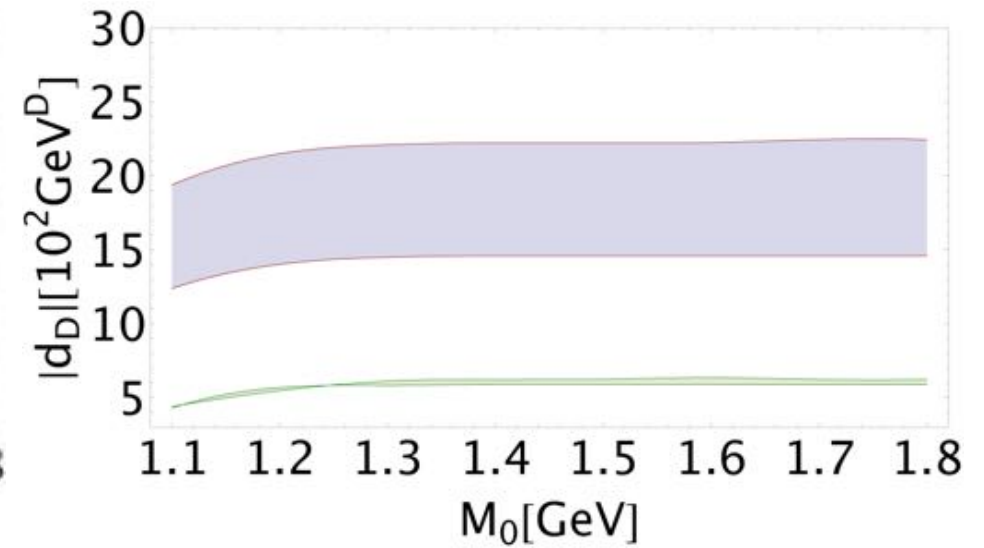
where (2) and (6) indicate the number of iterations.

Analysis from τ -like decay moments

● d_8 from R_1 : inputs α_s, d_4, d_{10}

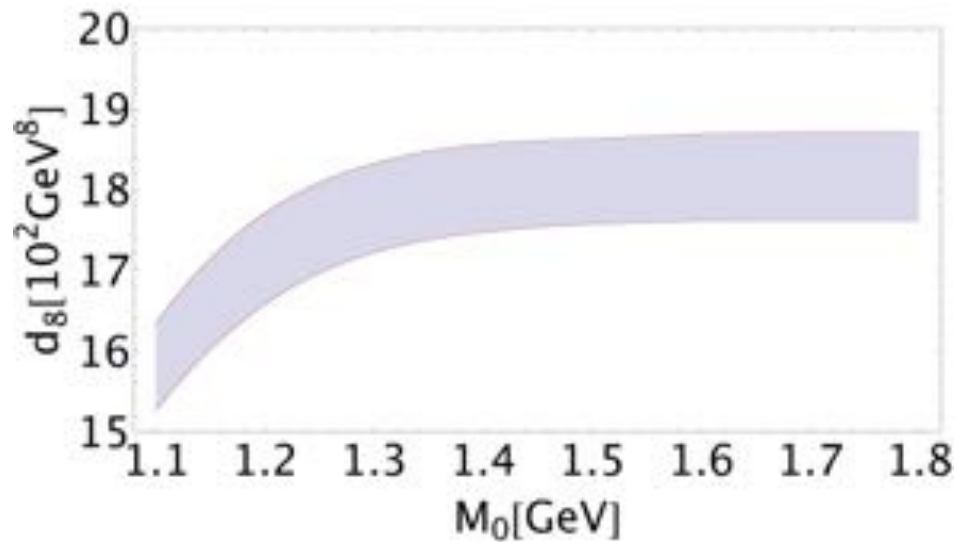


d_{10}, d_{12} from R_2 : inputs α_s, d_6

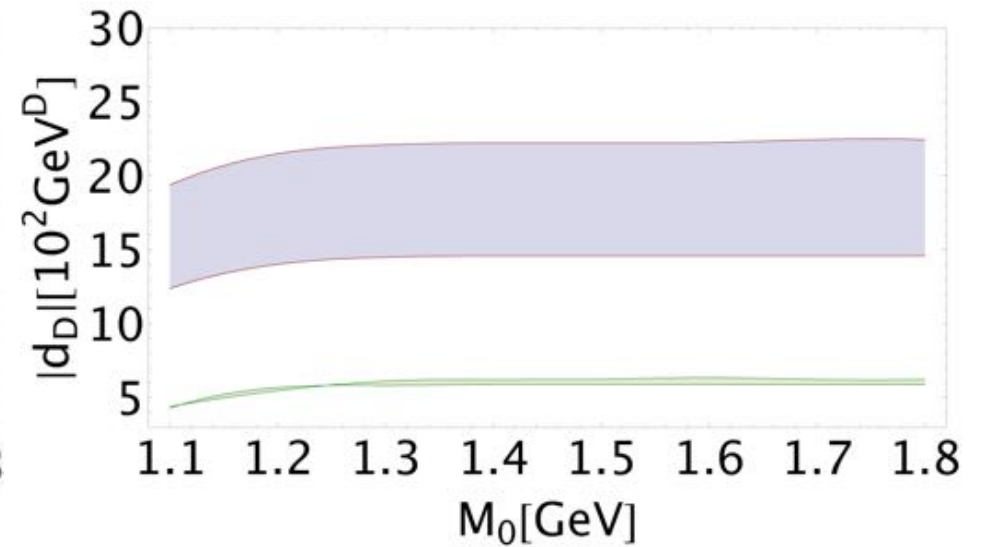


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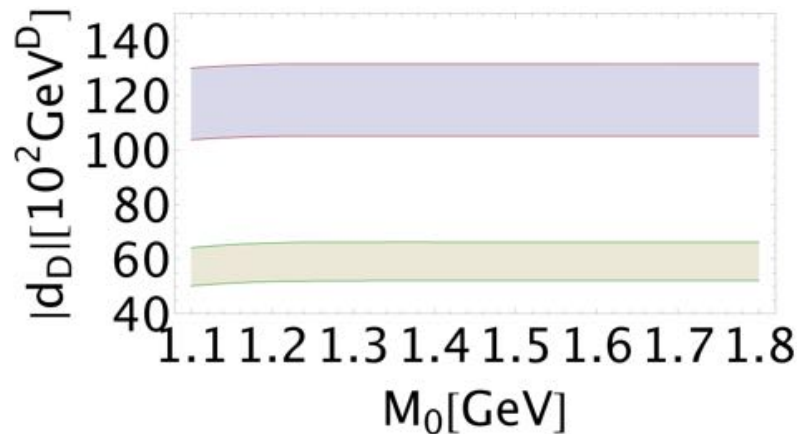
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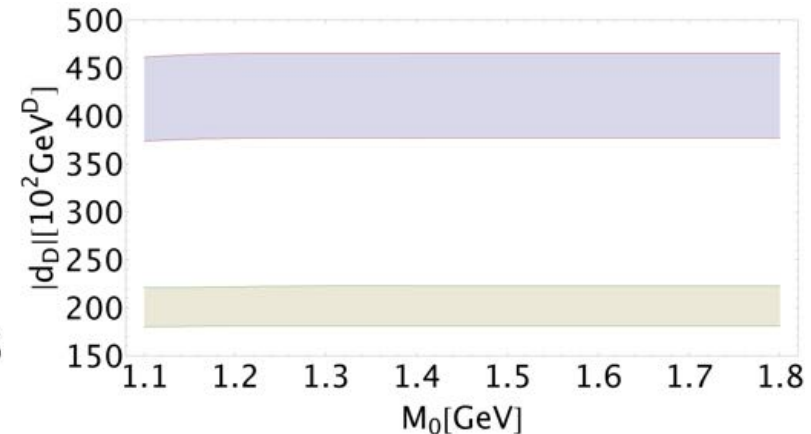
d_{10}, d_{12} from R_2 : inputs α_s, d_6



• d_{14}, d_{16} from R_3 and R_4



d_{18}, d_{20} from R_5 and R_6 .

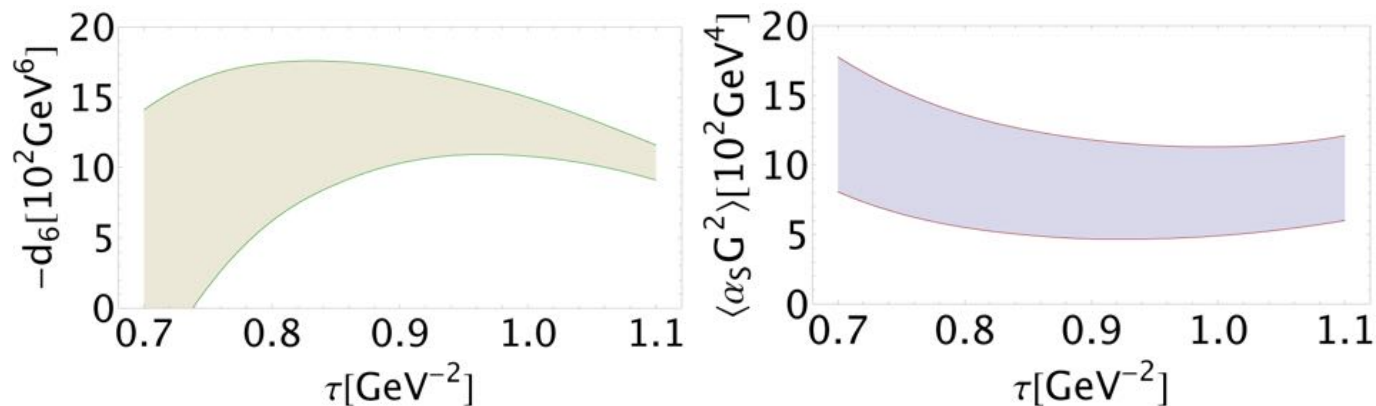


Re-analysis of the ratio of LSR $R_{10}(\tau)$

- Include the condensates of dimension $D \leq 16$ in the OPE.

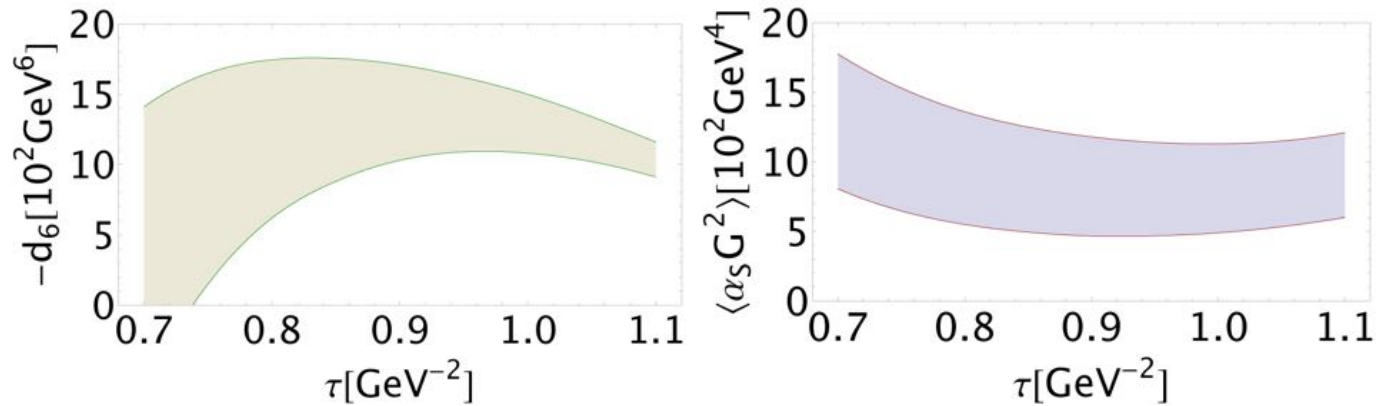
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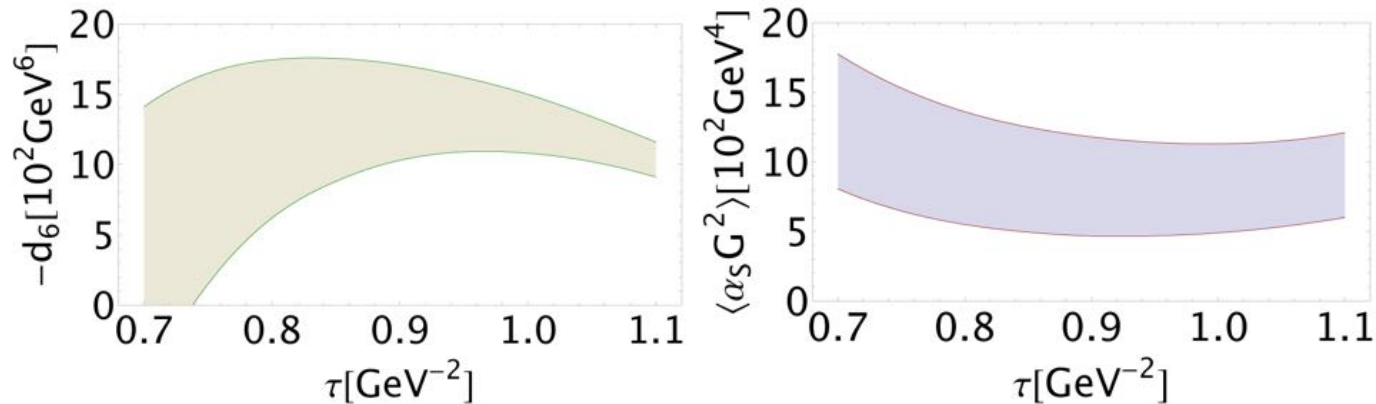
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- $d_6 = -(23.6 \pm 3.7) \times 10^{-2} \text{ GeV}^6$: for $\tau \simeq (0.85 \sim 0.95) \text{ GeV}^{-2}$
 - Agrees with $d_6 = -(20.5 \pm 2.2) \times 10^{-2} \text{ GeV}^6$ (initial inputs obtained at larger τ -values).
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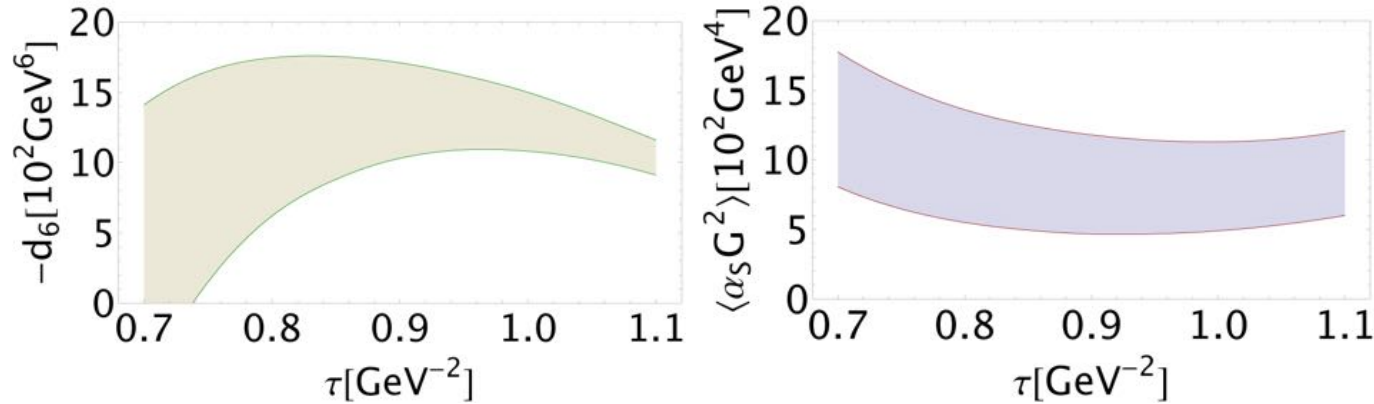
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- $\langle \alpha_s G^2 \rangle = (7.8 \pm 3.3) \times 10^{-2} \text{ GeV}^4$: for $\tau \simeq (0.9 \sim 1.0) \text{ GeV}^{-2}$
 - Agrees with $(6.49 \pm 0.35) \times 10^{-2} \text{ GeV}^4$ from Heavy quarks but less accurate !

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- Good convergence of the OPE at the optimization scale :

$$\begin{aligned} \tau R_{10}(\tau) &= 1 + \beta_1 a_s^2 + \dots - 0.123 \tau^2 + 0.164 \tau^3 - .140 \tau^4 - 0.015 \tau^5 + 0.050 \tau^6 - 0.034 \tau^7 + 0.193 \tau^8 \\ &= 1 + 10^{-3} + \sum_{n=2}^3 d_{2n} \tau^n = 0.127(0.091) + \sum_{n=4}^8 d_{2n} \tau^n = 0.017(-0.023) \text{ for } \tau = 0.95 (0.85) \text{ GeV}^{-2} \end{aligned}$$

Summary : QCD condensates from $e^+ e^- \rightarrow \text{Hadrons}$

- This work at (FO) : $\langle \alpha_s G^2 \rangle = (6.49 \pm 0.35) \times 10^{-2} \text{ GeV}^4$ input from $J/\psi, \Upsilon$ for $d_n \geq 6$. Units : 10^{-2} GeV^D : D dimension of the condensates. Systematic Errors not Included.

$\langle \alpha_s G^2 \rangle$	$-d_6$	d_8	$-d_{10}$	$-d_{12}$	d_{14}	$-d_{16}$	$-d_{18}$	d_{20}
7.8 ± 3.5	23.6 ± 3.7	18.2 ± 0.6	6.0 ± 0.2	18.4 ± 3.8	59.2 ± 7.1	118.4 ± 13.2	202.0 ± 20.7	421.3 ± 44.3

- $\langle \alpha_s G^2 \rangle$ confirms heavy and light quark SR results but less accurate.
- d_6 confirms a huge violation of factorization: factor (7 ± 1) LNT 84, SN96,
- d_8 confirms SN96
- $|d_{2n+2}| \approx 2|d_{2n}|$ for $n \geq 7$ (No Factorial / Exponential Growth !)
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- QCD condensates from some other τ -like decay moments at (FO)

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0.67 ± 0.89	15.2 ± 2.2	22.3 ± 2.5			ALEPH 99
5.34 ± 3.64	14.2 ± 3.5	21.3 ± 2.5			OPAL 99
$3.5_{-3.8}^{+2.2}$	$19.7_{-7.9}^{+11.8}$	$23.7_{-15.8}^{+11.8}$	11.8 ± 19.7	7.9 ± 19.7	Pich-Rodriguez 16 ($d_{14,16} = 0$)

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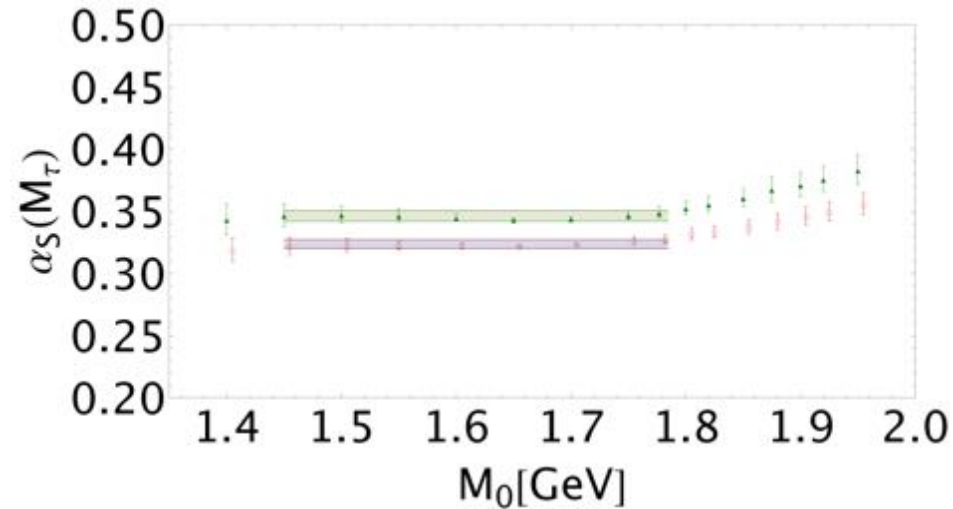
QCD COUPLING α_s
from $e^+e^- \rightarrow$ Hadrons
and $\tau \rightarrow$ Vector + ν_τ

$\alpha_s(M_\tau)$ *from τ -like decay BNP lowest moment*

- Use as inputs the previous condensates of dimension $d_4 \rightarrow d_8$.

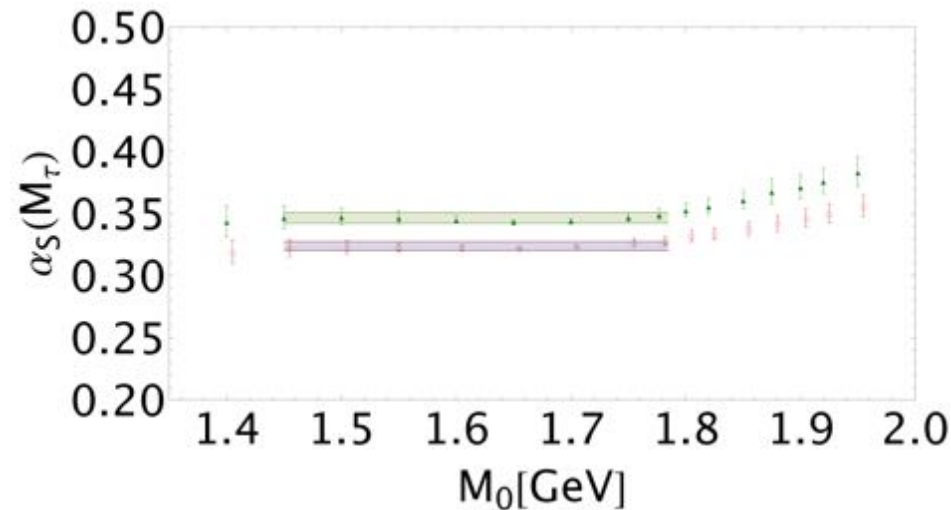
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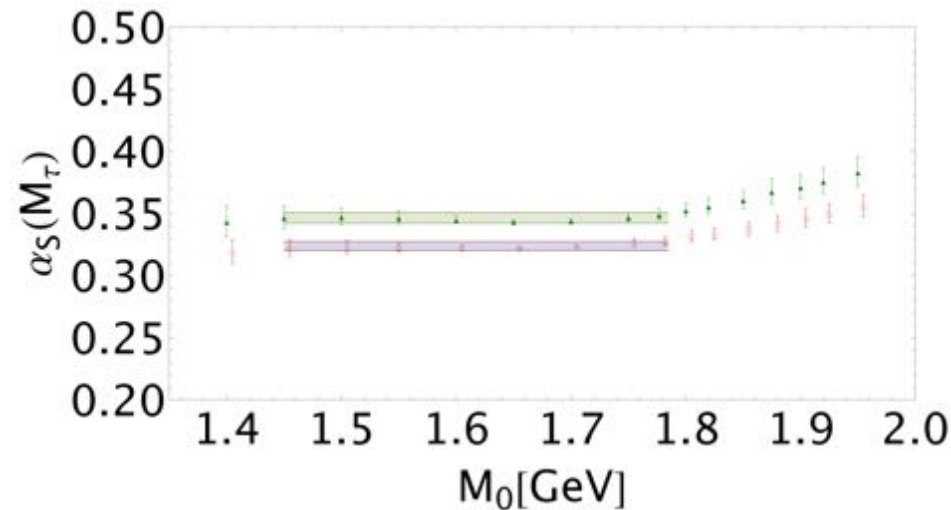


- One obtains to order $O(\alpha_s^4)$ [$M_0 = 1.675(25)$ GeV (stability point)] :

$$\begin{aligned}\alpha_s(M_\tau) &= 0.3238(36) \implies \alpha_s(M_Z) = 0.1190(2) && \text{FO} \\ &= 0.3465(43) \implies \alpha_s(M_Z) = 0.1216(2) && \text{CI}\end{aligned}$$

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- Sum of NP terms :

$$\sum \delta_{NP}(M_\tau) = (3.8 \pm 0.8) \times 10^{-2} \quad \approx \alpha_s^3 \quad (FO), \quad \alpha_s^2 \quad (CI).$$

Estimate of the error due to α_s^5

- Geometric growth of different PT series SN-Zakharov 09 :

$$D(Q^2) = \sum_n a_s^n c_n : c_0 = c_1 = 1, c_2 = 1.656, c_3 = 6.37, c_4 = 49.09$$

$$\implies c_4 \approx c_3^2 \implies c_5 \simeq (c_3/c_2) c_4^2 \approx (228 \pm 114).$$

$$R_0 = \sum_n a_s^n (g_n + c_n) : g_n \text{ from RG - resummation}$$

(see e.g : Pich – Lediberder92, Kataev – Starshenko95)

$$g_5 = -780 (FO), \quad 0(CI)$$

$$\implies \Delta\alpha_s(M_\tau) = \pm 71 \times 10^{-4} (FO), \quad \pm 62 \times 10^{-4} (CI),$$

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- $\alpha_s(M_\tau)$ to order α_s^4 including error due to α_s^5

$$\alpha_s(M_\tau) = 0.3238(36)_{fit}(71)_{\alpha_s^5} \implies \alpha_s(M_Z) = 0.1190(9)(3)_{evol} \quad \text{FO}$$

$$= 0.3465(43)_{fit}(62)_{\alpha_s^5} \implies \alpha_s(M_Z) = 0.1216(8)(3)_{evol} \quad \text{CI}$$

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- Mean** : assume that $\Delta\alpha_s^5$ absorbs the \neq between **FO** and **CI** at H.O :

$$\alpha_s(M_\tau) = 0.3358(55)(107)_{syst} \implies \alpha_s(M_Z) = 0.1204(7)(3)_{evol}$$

where **syst** comes from the distance of the mean central value to the ones of **FO** or **CI**.

$\alpha_s(M_\tau)$ from $e^+e^- \rightarrow I=1$ Hadrons & $\tau \rightarrow$ Vector

• Comparison of \neq Determinations

	THIS WORK		ALEPH 98	OPAL 99	PR 16	ALEPH 14
	e^+e^-	τ -decay				
FO	0.3081(86)	0.3128(79)	0.3200(220)	0.3230(160)	0.3200(150)	–
CI	0.3260(78)	0.3291(70)	0.3400(230)	0.3470(220)	0.3370(200)	0.3460(110)
FO \oplus CI	0.3179(100)	0.3219(105)				
$\delta^{(NP)} \times 10^{-2}$	3.8(8)	3.7(1.0)	2.0(3)	3.6(4)	1.7(3)	2.0(3)

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CI	0.3260(78)	0.3291(70)	0.3400(230)	0.3470(220)	0.3370(200)	0.3460(110)
FO \oplus CI	0.3179(100)	0.3219(105)				
$\delta^{(NP)} \times 10^{-2}$	3.8(8)	3.7(1.0)	2.0(3)	3.6(4)	1.7(3)	2.0(3)

• Average of our result for FO \oplus CI :

$$\alpha_s(M_\tau)|_{e^+e^-}^{SVZ} = 0.3179(58)(81)_{syst} \longrightarrow \alpha_s(M_Z)|_{e^+e^-}^{SVZ} = 0.1182(12)(3)_{evol}$$

$$\alpha_s(M_\tau)|_{\tau,V} = 0.3219(52)(91)_{syst} \longrightarrow \alpha_s(M_Z)|_{\tau,V} = 0.1187(13)(3)_{evol}$$

$\alpha_s(M_\tau)$ from $e^+e^- \rightarrow I = 1$ Hadrons & $\tau \rightarrow$ Vector

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CI	0.3260(78)	0.3291(70)	0.3400(230)	0.3470(220)	0.3370(200)	0.3460(110)
FO \oplus CI	0.3179(100)	0.3219(105)				
$\delta^{(NP)} \times 10^{-2}$	3.8(8)	3.7(1.0)	2.0(3)	3.6(4)	1.7(3)	2.0(3)

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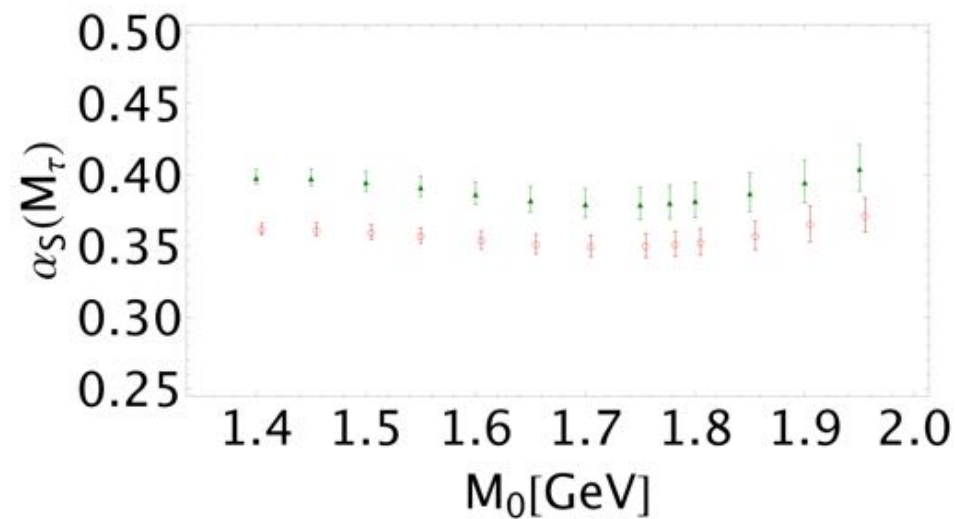
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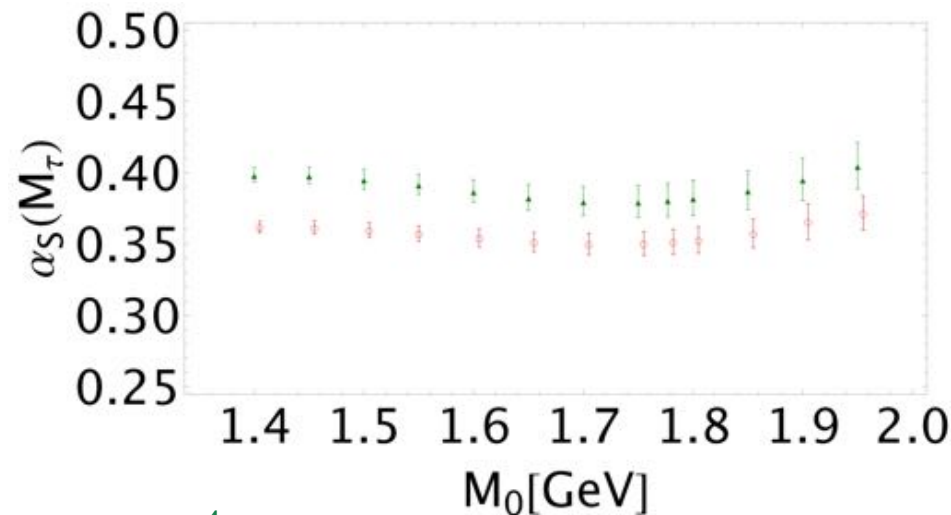
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- One obtains to order $O(\alpha_s^4)$:

$$\alpha_s(M_\tau) = 0.3514(68) \implies \alpha_s(M_Z) = 0.1221(7) \quad \text{FO}$$

$$= 0.3827(90) \implies \alpha_s(M_Z) = 0.1252(9) \quad \text{CI}$$

Relatively High compared to the PDG 24 average !

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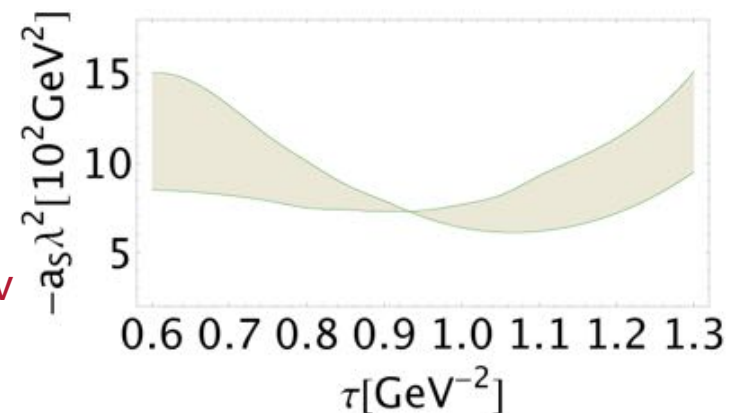
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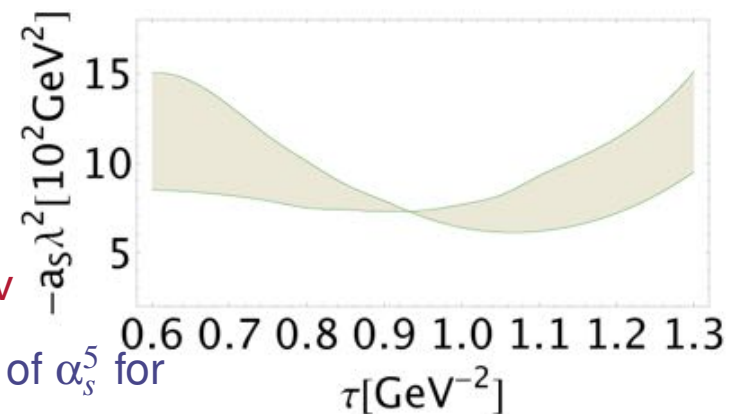
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- To avoid a double counting with the estimate of α_s^5 for

HO terms, we do not include by duality $\alpha_s \lambda^2$ but give an upper bound : $\delta_{\lambda^2}^{(2)} \leq 4\%$.



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- Contribute as operators of dimension 12 (running mass) NP 94:

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- Constraint from V-A τ -decay data and using : $\delta_{V-A}^{inst} \approx \delta_{V+A}^{inst}/20$ KK95:

$$\delta^{(9)}(M_\tau) \approx +(5 \sim 15) \times 10^{-4}$$

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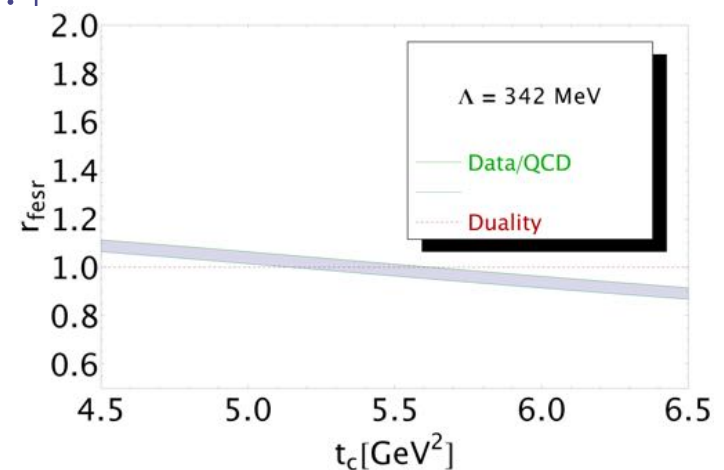
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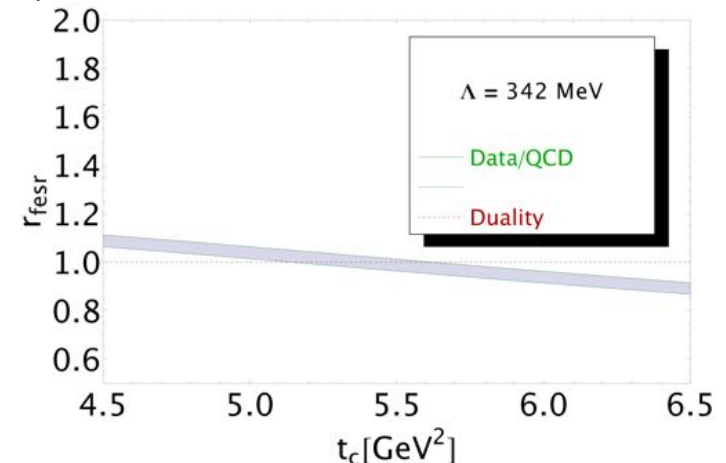
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SHIFMAN10.





Summary

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Misaotra Anareo @ ny Faharetana !

Merci pour Votre Patience !

Thanks for Your Patience !



<https://www.lupm.in2p3.fr/users/qcd/agmm>